Introduction

Katz Sándor

ELTE Elméleti Fizikai Tanszék
MTA-ELTE "Lendület" rácstérelmélet kutatócsoport

(Borsányi Szabolcs, Stephan Dürr, Fodor Zoltán, Christian Hoelbling, Stefan Krieg, Laurent Lellouch, Thomas Lippert, Antonin Portelli, Szabó Kálmán, Tóth Bálint)

SZTE Elméleti Fizikai Tanszék, 2016. február 18.







- Introduction
- Contributions to the proton/neutron masses
- Isospin symmetric case
- Isospin splittings
- Summary

Motivation

Introduction

Composition of the Universe

68% dark energy

27% dark matter

5% ordinary matter

Dark energy and dark matter are yet unknown

Ordinary matter (visible Universe):

Standard Model particles and their bound states mostly (in mass) atoms and nuclei 99.9% of the mass comes from protons & neutrons proton, neutron:

not elementary particles, built from quarks their masses are non-trivial, a consequence of the strong, electromagnetic and weak interactions

Isospin splittings

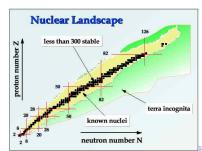
Summary

$$\Delta m_N = m_n - m_p = 0.0014 m_p$$

 Δm_N too small \rightarrow inverse β decay leaving predominantly neutrons $\Delta m_N \approx 0.05\%$ would already lead to much more He and much less H → stars would not have ignited as they did

 $\Delta m_N > 0.14\% \rightarrow$ much faster beta decay, less neutrons after BBN burning of H in stars and synthesis of heavy elements difficult

The whole nuclei chart is based on precise value of Δm_N





The 3 interactions in the Standard Model

```
Strong (QCD)
  guarks' internal degree of freedom: color
  SU(3) symmetry in color space
  coupling \approx 1
Electromagnetic
  among charged particles
  U(1) symmetry after the electroweak breaking
  coupling: \approx 1/137
Weak interaction
  original SU(2) \times U(1) symmetry spontaneously broken by the Higgs
  no exact symmetry at low energies
  approximate isospin symmetry is a remnant
  coupling: \approx 10^{-5}
+1: gravity
  coupling: \approx 10^{-40}
```

Origin of mass

Masses of elementary particles

Higgs mechanism:

Higgs-fermion Yukawa couplings (λ_i) lead to masses

$$m_i = \lambda_i v / \sqrt{(2)}$$

This is the contribution of the weak interaction

$$m_e = 0.511 \text{ MeV}$$

Quark masses are not uniquely defined. E.g. in $\overline{MS}(2\text{GeV})$ scheme:

$$m_u = 2.2(1) \text{ MeV}$$

$$m_d = 4.8(2) \text{ MeV}$$

Mass of the H atom: $m_H = 940 \text{ MeV}$

Where does the rest come from?



Mass of the proton and neutron

Proton and neutron are not elementary:

```
p = uud \quad n = udd
```

Besides the quark masses all energy (strong and EM) contributes

Strong interaction

quarks and gluons confined into the hadrons have huge kinetic and interaction energies This is the dominant contribution (99%)

Electromagnetic interaction

Charged particles have a static electric field which stores energy

Weak interaction

responsible for the guark masses through the Higgs mechanism negligible beyond that

Approximate solution

Isospin symmetry:

All properties of the u, d quarks are identical

their charges are neglected

good approximate symmetry but leads to identical m_n and m_n

Only QCD and the quark masses contribute

More detailed analysis

isospin symmetry is not exact

$$m_u \neq m_d$$
 and $Q_u = 2/3$ $Q_d = -1/3$

QCD + QED + quark masses contribute estimate:

quark masses: $\Delta m_N \approx 2.6 \text{ MeV}$

electrostatic energy of the proton: $\Delta m_N \approx -0.9 \text{ MeV}$



Quantum Chromodynamics (QCD)

QCD: Currently the best known theory to describe the strong interaction.

SU(3) gauge theory with fermions in fundamental representation.

Fundamental degrees of freedom:

• gluons:
$$A^a_{\mu}$$
, $a = 1, ..., 8$

• quarks: ψ , $3(color) \times 4(spin) \times 6(flavor)$ components

$$\mathcal{L}_{\text{QCD}} = \underbrace{-\frac{1}{4}F^{a}_{\mu\nu}F^{a\mu\nu}}_{\text{pure gauge part}} + \underbrace{\overline{\psi}(\textit{iD}_{\mu}\gamma^{\mu} - \textit{m})\psi}_{\text{fermionic part}},$$

where

$$F^a_{\mu\nu} = \partial_\mu A^a_
u - \partial_
u A^a_
u + g f^{abc} A^b_
u A^c_
u$$
 field strength

$$D_{\mu} = \partial_{\mu} + gA_{\mu}^{a} \frac{\lambda^{a}}{2i}$$
 covariant derivative \longrightarrow gives quark–gluon interaction

 \mathcal{L}_{QCD} is invariant under local gauge transformations:

$$A'_{\mu}(x) = G(x)A_{\mu}(x)G(x)^{\dagger} - \frac{i}{g} \left(\partial_{\mu}G(x)\right)G(x)^{\dagger}$$
 $\psi'(x) = G(x)\psi(x)$
 $\overline{\psi}'(x) = \overline{\psi}(x)G^{\dagger}(x)$

Only gauge invariant quantities are physical.

Properties of QCD:

Introduction

- Asymptotic freedom: Coupling constant $q \to 0$ when energy scale $\mu \to \infty$.
 - ⇒ Perturbation theory can be used at high energies.
- Confinement:
 - Coupling constant is large at low energies.
 - ⇒ Nonperturbative methods are required. → Nonperturbative methods are required.

Quantization using Feynman path integral:

$$\left\langle 0 | \ T[\mathcal{O}_1(x_1) \cdots \mathcal{O}_n(x_n)] \, | 0 \right\rangle = \frac{\int [\mathrm{d} \psi] \, [\mathrm{d} \overline{\psi}] \, [\mathrm{d} A_\mu] \, \mathcal{O}_1(x_1) \cdots \mathcal{O}_n(x_n) \, e^{iS\left[\psi, \overline{\psi}, A_\mu\right]}}{\int [\mathrm{d} \psi] \, [\mathrm{d} \overline{\psi}] \, [\mathrm{d} A_\mu] \, e^{iS\left[\psi, \overline{\psi}, A_\mu\right]}}$$

 e^{iS} oscillates \longrightarrow hard to evaluate integrals.

Wick rotation: $t \rightarrow -it$ analytic continuation to Euclidean spacetime.

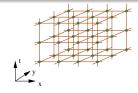
$$\implies$$
 $e^{iS} \longrightarrow e^{-S_E}$, where

$$S_{\mathsf{E}} = \int\!\mathrm{d}^4 x \,\, \mathcal{L}_{\mathsf{E}} = \int\!\mathrm{d}^4 x \, \left[rac{1}{4} F^a_{\mu
u} F^a_{\mu
u} + \overline{\psi} (D_\mu \gamma^\mu + m) \psi
ight]$$

positive definite Euclidean action.



Lattice regularization



take a finite spacetime volume (TV) and discretize with a lattice gluons live on links, quarks live on sites path integral \rightarrow finite ($\mathcal{O}(10^9)$) number of integrals quarks integrated analytically: $S \rightarrow S_{eff}$ numerical Monte Carlo integration possible for the gauge fields

Importance sampling

generate configurations with $\exp(-S_{eff})$ weight

Particle masses

two point functions: $\langle O(0)O(t)\rangle \longrightarrow \text{const} \cdot \exp(-\text{mt})$

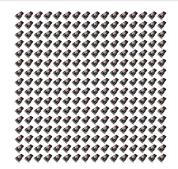


Juqueen system in Jülich

- 458752 cores
- 6 Pflop/s peak performance
- 2 Pflop/s sustained performance for Monte-Carlo codes







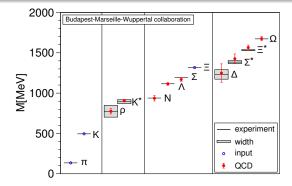
GPU cluster at ELTE

- 387072 cores

Introduction

- 1.1 Pflop/s peak performance
- 78 Tflop/s sustained performance





Science 322 (2008) 1224

masses extracted from various two-point functions correlation length → hadron masses perfect agreement with experiments



Unprecedented precision is required

 $\Delta m_N/m_N=0.14\%
ightarrow$ sub-permil precision is needed to get a high significance on Δm_N

 $m_u \neq m_d \rightarrow 1+1+1+1$ flavor lattice calculations are needed \rightarrow algorithmic challenge (Previous QCD calculations were typically 2+1 or 2+1+1 flavors)

Inclusion of QED: no mass gap

- → power-like finite volume corrections expected
- → long range photon field may cause large autocorrelations

Introduction

- Interpolation/extrapolation to physical point (continuum limit, physical quark masses)
- Treat power-like finite volume corrections analytically/numerically
- Potential precision: $\mathcal{O}(\frac{1}{N_c m_b^2}, \alpha^2) \sim 10^{-4}$
- Calculate mass splittings of N, Σ, Ξ as well as charmed hadrons, D, Ξ_{cc} and Coleman-Glashow relation
- Reaching the desired accuracy via a blind analysis (hard)
- Determine QCD/QED contributions of splittings



Isospin splittings

Summary

Zero mode subtraction

The absence of a mass gap may cause divergences at finite volume perturbative momentum sums $\rightarrow 1/k^2$ factors \rightarrow zero mode problematic

Removing a finite number of modes does not change $V \to \infty$ physics

Advantages of zero mode removal:

- \rightarrow analytic calculation of finite V corrections is possible
- → algorithmic speedup

Many possibilities, we use:

(Hayakawa, Uno) all spatial zero modes:
$$\sum_{ec{x}} A_{\mu,x_0,ec{x}} = 0 \quad orall \mu, x_0$$



QED in a finite volume

Calculate 1 loop self energy of charged particles in finite and infinite volumes

For a point-like particle

$$\Delta\Sigma(p,L) = \left[\sum_{k} - \int \frac{d^{4}k}{(2\pi)^{4}}\right] \sigma(k,p)$$

zero mode removal: $\sum_{\nu} \equiv \frac{1}{TL^3} \sum_{k_0} \sum_{\vec{k} \neq 0}$

Finite V correction to the pole mass can be calculated Result for a spin half particle:

$$m(L) \underset{L \to +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[\frac{\kappa}{2mL} \left(1 + \frac{2}{mL} \right) - \frac{3\pi}{(mL)^3} \right] \right\}$$

with $\kappa = 2.837297(1)$



In QCD charged hadrons are not point-like previous results have to be extended

QED Ward identities \rightarrow first two orders universal:

$$m(T,L) \underset{T,L \to +\infty}{\sim} m \left\{ 1 - q^2 \alpha \frac{\kappa}{2mL} \left[1 + \frac{2}{mL} \right] + \mathcal{O}(\frac{\alpha}{L^3}) \right\}$$

for scalars and spin half fermions

Form factors (e.g. charge radius) enter at $\mathcal{O}(\frac{\alpha}{13})$ level.

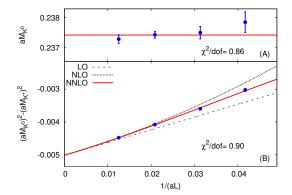
Strategy:

include analytic corrections for the two universal orders fit coefficient of 1/L3

1/L3 in many cases negligible, only significant for mesons



Finite V dependence of the kaon mass



Neutral kaon shows no volume dependence Volume dependence of the K splitting is perfectly described 1/L³ order is significant



Determining the isospin splittings

two sources of isospin violation: electromagnetism & $m_{IJ} \neq m_{cl}$ we work at various elecromagnetic couplings and quark masses renormalized coupling defined at hadronic scales

 \Rightarrow only linear term in α

 $\delta m = m_d - m_u$ is very small \Rightarrow linear term is also enough for δm

$$\Rightarrow \Delta M_X = F_X(M_{\pi^+}, M_{K^0}, M_{D^0}, L, a) \cdot \alpha + G_X(M_{\pi^+}, M_{K^0}, M_{D^0}, a) \cdot \Delta M_K^2$$

 ΔM_K^2 : QED-like L dependence: absorbed in $F_X(..., L, a)$ charged particle masses: corrected for universal finite-size effect non-universal effects starting with $1/L^3$ are allowed in the QED part alternative procedure: use ΔM_{Σ} (less precise) instead of ΔM_{Σ}^2

if α or $(m_{ij} - m_{ij})$ vanishes QED or QCD parts disappear separation: $\Delta M_X = \Delta_{\rm OED} M_X + \Delta_{\rm OCD} M_X$ it is sufficient to decompose the kaon mass squared difference use $\Delta M_X = F_X(M_{\pi^+}, M_{K^0}, M_{D^0}, L, a) \cdot \alpha + G_X(M_{\pi^+}, M_{K^0}, M_{D^0}, a) \cdot \Delta M_K^2$ separation ambigous: depends on the choice of scheme for $m_{II} - m_{dI}$

Suggested separation:

use Σ^+ and Σ^- baryons

they have the same charge squared and same spin

leading order: mass difference comes from quark masses only

Introduction

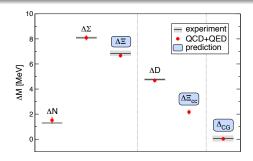
Summary

suggestion: $\Delta_{CG} = \Delta M_N - \Delta M_{\Sigma} + \Delta M_{\Xi}$ is close to zero remember the quark compositions, in which it cancels indeed: M(ddu) - M(uud) - M(sdd) + M(suu) + M(ssd) - M(ssu) = 0determine the leading order terms in the α and δm expansion for $\alpha=0$ a complete quark exchange symmetry $\Delta_{\rm CG} \propto (m_{\rm s}-m_{\rm d})(m_{\rm s}-m_{\rm u})(m_{\rm d}-m_{\rm u})$ for $\alpha > 0$ remains a $d \leftrightarrow s$ symmetry, thus $\Delta_{CG} \propto \alpha (m_s - m_d)$

the Coleman-Glashow relation is satisfied to high accuracy $\Delta_{\text{CG}} = 0.00(11)(06) \text{ MeV}$



Introduction



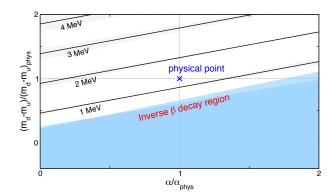
Science 347 (2015) 1452

	splitting [MeV]	QCD [MeV]	QED [MeV]
ΔN=n-p	1.51(16)(23)	2.52(17)(24)	-1.00(07)(14)
$\Delta \Sigma = \Sigma^{-} - \Sigma^{+}$	8.09(16)(11)	8.09(16)(11)	0
ΔΞ=ΞΞ0	6.66(11)(09)	5.53(17)(17)	1.14(16)(09)
$\Delta D=D^{\pm}-D^{0}$	4.68(10)(13)	2.54(08)(10)	2.14(11)(07)
$\Delta \Xi_{cc} = \Xi_{cc}^{++} - \Xi_{cc}^{+}$	2.16(11)(17)	-2.53(11)(06)	4.69(10)(17)
$\Delta_{\text{CG}} = \Delta N - \Delta \Sigma + \Delta \Xi$	0.00(11)(06)	-0.00(13)(05)	0.00(06)(02)

Summary

Dependence on fundamental parameters

How strongly does Δm_N depend on α and the guark masses?



- Mass of the visible universe: 99.9% protons and neutrons
- Their masses almost identical, but the tiny difference has cosmological consequences
- Higgs is responsible for 1% of the mass
- The rest is given mostly by QCD and a small contribution by QED
- Both masses have been determined using lattice calculations
- The mass difference is a result of two competing effects (quark masses and charges)
- Good agreement with experiments + QCD/QED contributions given separately
- 4 further splittings determined, 2 of them are predictions

