

Exceptional orthogonal polynomials and quasi-invariance

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Bochner's theorem (1929)

Suppose $p_n(x) \in \mathbb{R}[x]$, $n \in \mathbb{Z}_{\geq 0}$, with $\deg p_n = n$ satisfy

$$A_2(x) \frac{d^2 p_n}{dx^2} + A_1(x) \frac{dp_n}{dx} + A_0(x) p_n = E_n p_n(x).$$

Then $A_j(x) \in \mathbb{R}[x]$ with $\deg A_j \leq j$.

If, in addition,

$$\int_a^b p_m(x) p_n(x) w(x) dx = \delta_{mn} g_n,$$

with $w(x) > 0$, then (up to $x \rightarrow ax + b$)
the $p_n(x)$ are classical Hermite, Laguerre or
Jacobi polynomials.



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Hermite polynomials

The classical Hermite polynomials

$$H_n(z) = (-1)^n e^{z^2} \frac{d^n}{dz^n} e^{-z^2} = n! \sum_{m=0}^{\lfloor n/2 \rfloor} \frac{(-1)^m}{m!(n-2m)!} (2z)^{n-2m}$$

satisfy

$$\frac{d^2 H_n}{dz^2} - 2z \frac{dH_n}{dz} = -2nH_n$$

and

$$\int_{-\infty}^{\infty} H_n(z) H_m(z) e^{-z^2} dz = \delta_{nm} 2^n n! \sqrt{\pi}.$$

Exceptional orthogonal polynomials

Let $S \subset \mathbb{Z}_{\geq 0}$ be such that $|\mathbb{Z}_{\geq 0} \setminus S| < \infty$.

Gómez-Ullate, Kamran and Milson (2010) called $p_n(x) \in \mathbb{R}[x]$, $n \in S$, a system of **exceptional orthogonal polynomials** if the following conditions are satisfied.

Eigenvalue equation:

$$A_2(x) \frac{d^2 p_n}{dx^2} + A_1(x) \frac{dp_n}{dx} + A_0(x) p_n = E_n p_n(x).$$

Orthogonality:

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Density: $U := \langle p_n : n \in S \rangle$ is dense in $\mathbb{R}[x]$, i.e.

$$(p, p_n) = 0 \quad \forall n \in S \quad \Rightarrow \quad p \equiv 0.$$

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Exceptional Hermite polynomials

Let $\lambda = (\lambda_1, \dots, \lambda_l)$ be a double partition:

$$\lambda = \mu^2 = (\mu_1, \mu_1, \mu_2, \mu_2, \dots, \mu_r, \mu_r), \quad 2r = l,$$

and

$$k_j = \lambda_j + l - j, \quad j = 1, \dots, l.$$

Gómez-Ullate, Kamran and Milson (2014) showed that

$$H_{\lambda,n}(z) := \text{Wr}(H_n, H_{k_1}, \dots, H_{k_l}), \quad n \in \mathbb{Z}_{\geq 0} \setminus \{k_1, \dots, k_l\},$$

yield a system of **exceptional orthogonal polynomials**.

In particular,

$$\int_{-\infty}^{\infty} H_{\lambda,n}(z) H_{\lambda,m}(z) \frac{e^{-z^2}}{W_{\lambda}(z)^2} = \delta_{nm} \sqrt{\pi} 2^n n! \prod_{i=1}^l (n - k_i)(n - k_i - 1),$$

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Singular weight functions

Why restrict attention to double partitions?

Crum (1954) and Adler (1994): The Wronskian $W_\lambda(z)$ has no zeros on \mathbb{R} iff

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Our aim: Obtain a **natural interpretation** of the polynomials

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The simplest non-trivial example: $\lambda = (1)$

The first few polynomials $H_{(1),n} = \text{Wr}(H_n, H_1)$ are

$$\begin{aligned} H_{(1),0}(z) &= 2 && \text{degree 1 missing} \\ H_{(1),2}(z) &= -(2 + 4z^2) && H_{(1),3}(z) = -16z^3 \\ H_{(1),4}(z) &= 12(1 + 4z^2 - 4z^4) && H_{(1),5}(z) = 64z^3(5 - 2z^2) \end{aligned}$$

Eigenvalue equation:

$$\frac{d^2 H_{(1),n}}{dz^2} - 2 \left(z + \frac{1}{z} \right) \frac{dH_{(1),n}}{dz} = -2(1+n)H_{(1),n}.$$

Orthogonality:

$$\int_{i\xi + \mathbb{R}} H_{(1),n}(z) H_{(1),m}(z) \frac{e^{-z^2}}{(2z)^2} dz = \delta_{nm} 2^n n! \sqrt{\pi} 2(n-1), \quad \xi \neq 0.$$

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Darboux transformations

Harmonic oscillator:

$$\mathcal{H} := -\frac{d^2}{dz^2} + z^2$$

has eigenfunctions

$$\psi_n(z) = H_n(z)e^{-z^2}, \quad n \in \mathbb{Z}_{\geq 0},$$

w/ eigenvalues $E_n = 2n + 1$.

Letting

$$D_k = \frac{d}{dz} - \frac{\psi'_k}{\psi_k},$$

we get

$$\begin{aligned} D_k^* D_k &= \left(-\frac{d}{dz} - \frac{\psi'_k}{\psi_k} \right) \left(\frac{d}{dz} - \frac{\psi'_k}{\psi_k} \right) \\ &= -\frac{d^2}{dz^2} + \frac{\psi''_k}{\psi_k} \\ &= -\frac{d^2}{dz^2} + z^2 - E_k, \end{aligned}$$

so that

$$\mathcal{H} = D_k^* D_k + E_k.$$

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gives

$$\mathcal{H}_k := D_k D_k^* + E_k = -\frac{d^2}{dz^2} + z^2 - 2 \left(\frac{\psi'_k}{\psi_k} \right)'.$$

Introducing

$$\psi_{k,n} := D_k \psi_n = \frac{\text{Wr}(\psi_n, \psi_k)}{\psi_k}, \quad n \neq k,$$

we obtain

$$\mathcal{H}_k \psi_{k,n} = D_k (D_k^* D_k + E_k) \psi_n = D_k \mathcal{H} \psi_n = E_n \psi_{k,n}.$$

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In particular,

$$\mathcal{H}_1 = -\frac{d^2}{dz^2} + z^2 + \frac{2}{z^2}$$

and

$$\begin{aligned}\psi_{1,n}(z) &= \frac{\text{Wr}(H_n(z)e^{-z^2/2}, H_1(z)e^{-z^2/2})}{H_1(z)e^{-z^2/2}} \\ &= \frac{\text{Wr}(H_n(z), H_1(z))}{H_1(z)} e^{-z^2/2} \\ &= H_{(1),n}(z) \frac{e^{-z^2/2}}{2z}.\end{aligned}$$

Orthogonality

$$\begin{aligned}\int_{i\xi+\mathbb{R}} H_{(1),n}(z)H_{(1),m}(z)\frac{e^{-z^2}}{(2z)^2}dz &= \int_{i\xi+\mathbb{R}} \psi_{1,n}(z)\psi_{1,m}(z)dz \\ &= \int_{i\xi+\mathbb{R}} \psi_n(z)(D_1^*D_1\psi_m(z))dz \\ &= \int_{i\xi+\mathbb{R}} \psi_n(z)([\mathcal{H} - E_1]\psi_m(z))dz \\ &= 2(m-1)\int_{i\xi+\mathbb{R}} \psi_n(z)\psi_m(z)dz \\ &= 2(m-1)\int_{\mathbb{R}} H_n(z)H_m(z)e^{-z^2}dz \\ &= \delta_{nm}2^n n! \sqrt{\pi} 2(n-1)\end{aligned}$$

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Quasi-invariance

Note that

$$\begin{aligned}\frac{d}{dz}(z\psi_{1,n}) &= \frac{d}{dz}(H'_n H_1 - H_n H'_1) \frac{e^{-z^2/2}}{2} \\ &= \frac{d}{dz}(zH'_n - H_n) e^{-z^2/2} \\ &= z(\dots) e^{-z^2/2} = 0 \quad \text{at } z = 0,\end{aligned}$$

i.e.

$$\psi_{1,n}(-z) = (-1)\psi_{1,n}(z) + o(z).$$

More generally, we say that a meromorphic function $\psi(z)$ is **quasi-invariant** at $z = z_i$ w/ multiplicity $m_i \in \mathbb{Z}_{\geq 0}$ if

1. $\psi(z)(z - z_i)^{m_i}$ is analytic at $z = z_i$,
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Let

$$\mathcal{Q}_{(1)} = \left\{ p \in \mathbb{C}[z] : \psi(z) = p(z) \frac{e^{-z^2/2}}{2z} \text{ q. - inv. at } z = 0 \text{ w/ mult. } 1 \right\}.$$

Proposition

We have

$$\mathbb{C}\langle H_{(1),n} : n \neq 1 \rangle = \mathcal{Q}_{(1)}.$$

Proof.

The subspaces have the same codimension in $\mathbb{C}[z]$, since

$$\# \text{ degrees missing} = 1 = \# \text{ q. - inv. conditions.}$$

□

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Quasi-invariance

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$$\mathcal{Q}_{(1)} = \left\{ p \in \mathbb{C}[z] : \psi(z) = p(z) \frac{e^{-z^2/2}}{2z} \text{ q. - inv. at } z = 0 \text{ w/ mult. } 1 \right\}.$$

Proposition

We have

$$\mathbb{C}\langle H_{(1),n} : n \neq 1 \rangle = \mathcal{Q}_{(1)}.$$

Proof.

The subspaces have the same codimension in $\mathbb{C}[z]$, since

$$\# \text{ degrees missing} = 1 = \# \text{ q. - inv. conditions.}$$

□

A Hermitian product

Definition

Assuming $\xi \in \mathbb{R} \setminus \{0\}$, we set

$$\langle p, q \rangle = \int_{i\xi + \mathbb{R}} p(z) \bar{q}(z) \frac{e^{-z^2}}{(2z)^2} dz, \quad p, q \in \mathcal{Q}_{(1)},$$

where $\bar{q}(z) = \overline{q(\bar{z})}$.

Lemma

$\langle \cdot, \cdot \rangle$ does not depend on the value of ξ .

Proof.

By the residue thm and quasi-invariance, we have

$$\begin{aligned} \langle p, q \rangle_{\xi} &= \langle p, q \rangle_{-\xi} + \frac{d}{dz} \left(z^2 p(z) \bar{q}(z) \frac{e^{-z^2}}{(2z)^2} \right) \Big|_{z=0} \\ &= \langle p, q \rangle_{-\xi} + \frac{d}{dz} \left(\left[z \cdot p(z) \frac{e^{-z^2/2}}{2z} \right] \cdot \left[z \cdot \bar{q}(z) \frac{e^{-z^2/2}}{2z} \right] \right) \Big|_{z=0} \\ &= \langle p, q \rangle_{-\xi}. \end{aligned}$$

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A Hermitian product

Proposition

$\langle \cdot, \cdot \rangle$ is Hermitian.

Proof.

Introducing

$$w(z) = \frac{e^{-z^2}}{(2z)^2}$$

and observing $\bar{w}(z) = w(z)$, we deduce

$$\begin{aligned}\langle p, q \rangle_\xi &= \int_{\mathbb{R}} p(i\xi + x) \bar{q}(i\xi + x) w(i\xi + x) dx \\ &= \overline{\int_{\mathbb{R}} \bar{p}(-i\xi + x) q(-i\xi + x) w(-i\xi + x) dx} \\ &= \overline{\langle q, p \rangle_{-\xi}}.\end{aligned}$$

Hence, the assertion follows from the lemma. □

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Main result

Theorem

The Exceptional Hermite Polynomials $H_{(1),n}$, $n \neq 1$, constitute an orthogonal basis in $\mathcal{Q}_{(1)}$ w/ respect to the Hermitian product $\langle \cdot, \cdot \rangle$. Specifically,

$$\langle H_{(1),n}, H_{(1),m} \rangle = \delta_{nm} 2^n n! \sqrt{\pi} 2(n-1).$$

This result generalises to Exceptional Hermite Polynomials labelled by an arbitrary partition λ , see

W. A. Haese-Hill, M. H. & A. P. Veselov (2016): Complex exceptional orthogonal polynomials and quasi-invariance, Lett. Math. Phys. 106.

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A multivariate generalisation

The quantum rational (A_{n-1}) Calogero-Moser system

Can be defined by the Hamiltonian

$$H_n = - \sum_{j=1}^n \frac{\partial^2}{\partial z_j^2} + 2m(m-1) \sum_{1 \leq j < k \leq n} \frac{1}{(z_j - z_k)^2},$$

where $N \in \mathbb{N}$ (particle number) and $m > 0$ (coupling parameter).

Associated **integrable system** (n commuting PDOs):

$$H_n^1 = -i \sum_{j=1}^n \frac{\partial}{\partial z_j}, \quad H_n,$$

$$H_n^k = (-i)^k \sum_{j=1}^n \frac{\partial^k}{\partial z_j^k} + \dots, \quad k = 3, \dots, n,$$

(*Olshanetsky & Perelomov (1983), Ujino, Hikami & Wadati (1992), and others*).

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A_{n-1} quasi-invariants

Let $m \in \mathbb{Z}_{\geq 0}$.

A polynomial $q \in \mathbb{C}[x]$ is said to be m -quasi-invariant if

$$\left(\frac{\partial}{\partial z_i} - \frac{\partial}{\partial z_j} \right) q(z) = \left(\frac{\partial}{\partial z_i} - \frac{\partial}{\partial z_j} \right)^3 q(z) = \cdots = \left(\frac{\partial}{\partial z_i} - \frac{\partial}{\partial z_j} \right)^{2m-1} q(z) \equiv 0$$

along

$$z_i = z_j, \quad \forall 1 \leq i < j \leq n.$$

Such polynomials form an algebra, denoted \mathcal{Q}_m .

The $m \in \mathbb{Z}_{\geq 0}$ Calogero-Moser system is algebraically integrable: for each $q \in \mathcal{Q}_m$ exists PDO

$$H_q = q(-i\partial/\partial z_1, \dots, -i\partial/\partial z_n) + \cdots$$

such that

$$H_{z_1^2 + \cdots + z_n^2} = H_n, \quad [H_q, H_p] = 0, \quad \forall q, p \in \mathcal{Q}_m,$$

(Chalykh & Veselov (1990)).

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Bilinear forms on \mathcal{Q}_m

For $q \in \mathcal{Q}_m$, let

$$L_q = A_m(z) \circ H_q \circ A_m(z)^{-1}, \quad A_m(z) = \prod_{1 \leq i < j \leq n} (z_i - z_j)^m.$$

Canonical:

$$\langle q, p \rangle := (L_q p)(0).$$

'Hermite-like':

$$(q, p) := \int_{i\xi + \mathbb{R}^n} \frac{q(z)p(z)}{A_m(z)^2} e^{-z^2} dz.$$

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A result in progress

Feigin, H. & Veselov: There exists a degree-preserving operator of the form

$$H_\bullet : \mathcal{Q}_m \rightarrow \mathcal{Q}_m, q \mapsto \langle F(z, \bullet), q \rangle$$

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and

$$\left(L_{z_1^2 + \dots + z_n^2} - 2 \sum_{i=1}^n z_i \frac{\partial}{\partial z_i} \right) H_q = -2(\deg q) H_q$$

whenever q is homogeneous.

The (generating) function F is given by

$$F(z, w) = \phi(z, w) \exp(-w^2),$$

w/ ϕ the so-called (A_{n-1}) rational Baker-Akhiezer function.

The result extends to all Coxeter groups.

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